

Advanced Simulation Methods for Charge Transport in OLEDs

Evelyne Knapp, B. Ruhstaller

Overview

1. Introduction
2. Physical Models
3. Numerical Methods
4. Outlook

www.icp.zhaw.ch

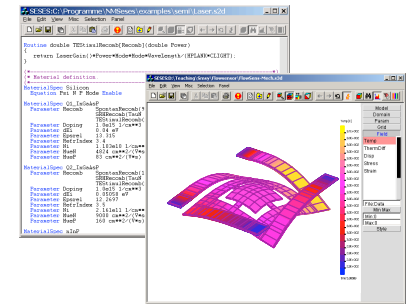
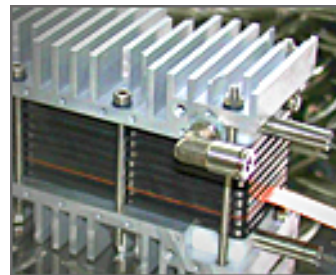
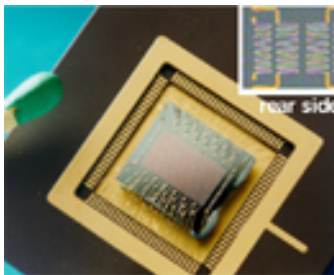
- Interdisciplinary team of 8 physicists, 4 mathematicians und 3 engineers



1996 Section NMSA
2002 Foundation CCP
2007 Foundation ICP

Spin-offs:
Numerical Modeling GmbH, www.nmtec.ch
Fluxim AG, www.fluxim.com

- The main focus is applied research and development in the following areas:
 - › Micro systems, sensors, actors
 - › Fuel cells
 - › Organic optoelectronic and photovoltaics
 - › Simulation software





Advanced Experimentally Validated OLED model

Philips Research Eindhoven

Project Coordinator:
Reinder Coehoorn

Philips Research Aachen

Zürich University of Applied Sciences

Fluxim

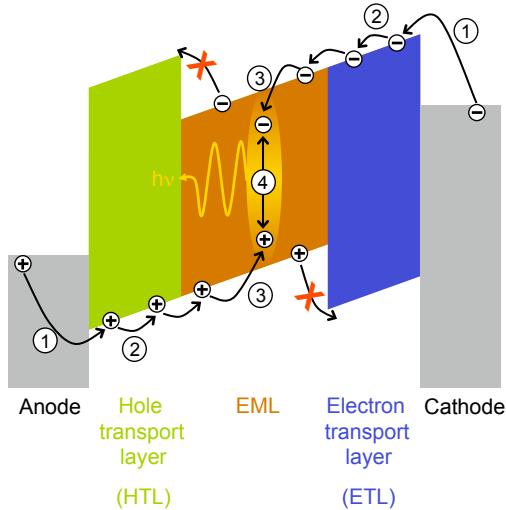
Eindhoven University of Technology

Technical University Dresden

Sim4Tec

University of Groningen

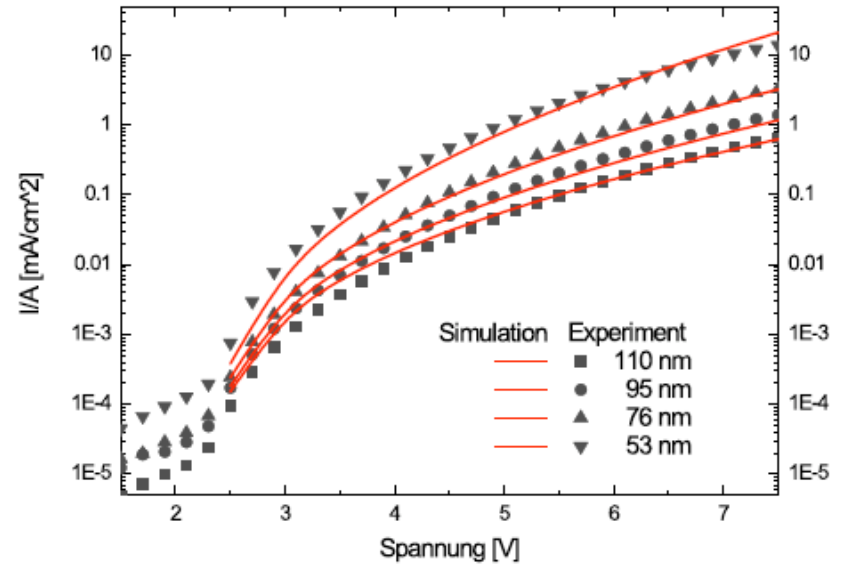
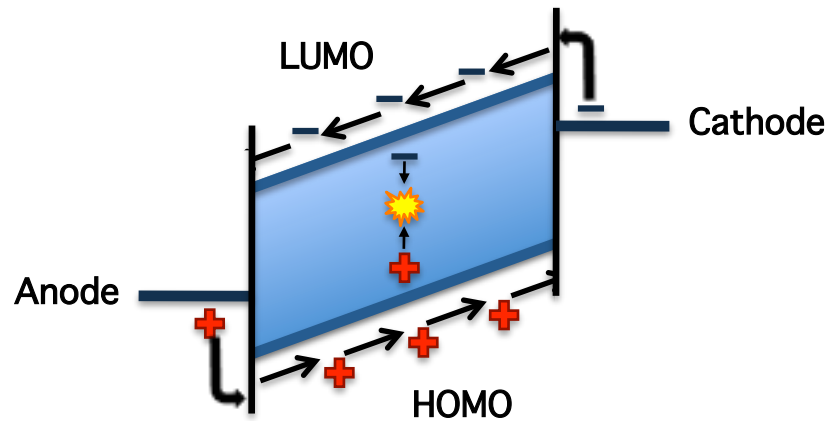
University of Cambridge



Fundamental Processes:

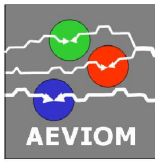
1. Charge Injection
2. Charge Carrier Transport
3. Exciton Formation
4. Radiative Decay
5. Light Extraction

Real stack consists of up to 12 layers!



- Novel physical models require better numerical methods
- Transient simulations and IV curves need multiple simulations
- ➔ Efficient simulations are crucial

Experimental data from CSEM, simulation by ICP

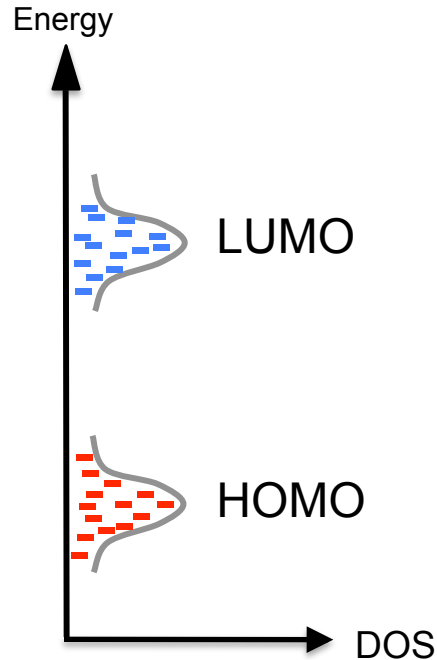


Overview-Task list



- ✓ Modeling of charge carrier transport
 - › Gummel solver
 - › Newton solver
- ✓ Bipolar
- ✓ Injection
- ✓ Organic material properties
 - › Disorder (Gaussian DOS)
 - › Mobility
 - › Generalized Einstein relation
- ✓ Traps (Exponential DOS)
- ✓ Multilayer OLEDs
 - Exciton dynamics
 - Parameter extraction
 - Coupling to optical model
 - Impedance simulations

Gaussian Disorder



- Small molecules and polymer LEDs/solar cells
- Charge transport by hopping between uncorrelated sites
- Width of DOS-disorder parameter σ (50-150 meV)



$$DOS(\epsilon) = \frac{N_t}{\sqrt{2\pi}\sigma} \exp \left[- \left(\frac{\epsilon - \epsilon_0}{\sqrt{2}\sigma} \right)^2 \right]$$

Poisson equation: $\epsilon \Delta \psi = q(n - p)$

Continuity equation: $\nabla \cdot J_p + q \frac{\partial p}{\partial t} = -qR(p, n)$

Drift-Diffusion: $J_p = -q\mu_p p \nabla \psi - qD_p \nabla p$

similar for electrons


Poisson equation: $\epsilon \Delta \psi = q(n - p)$

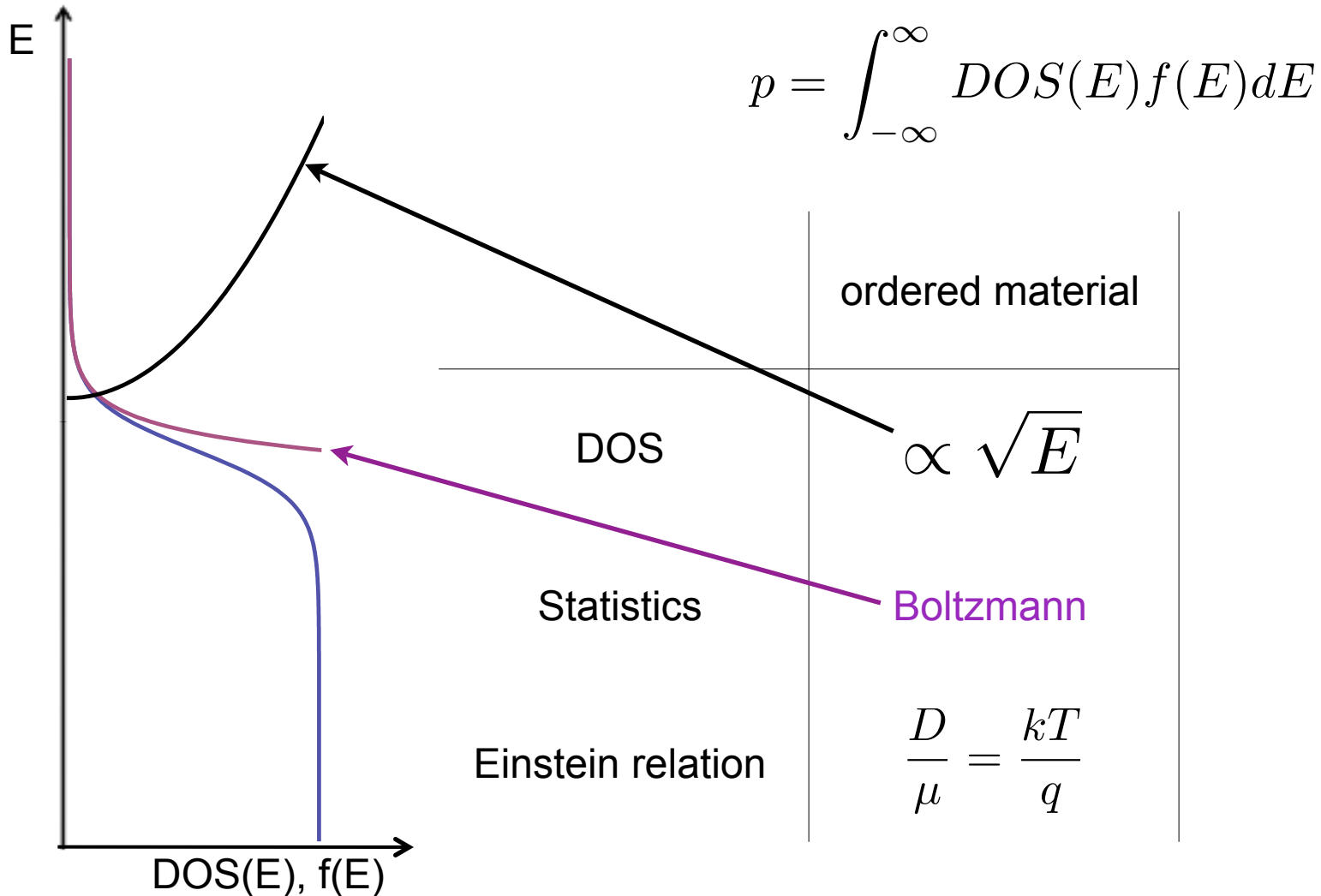
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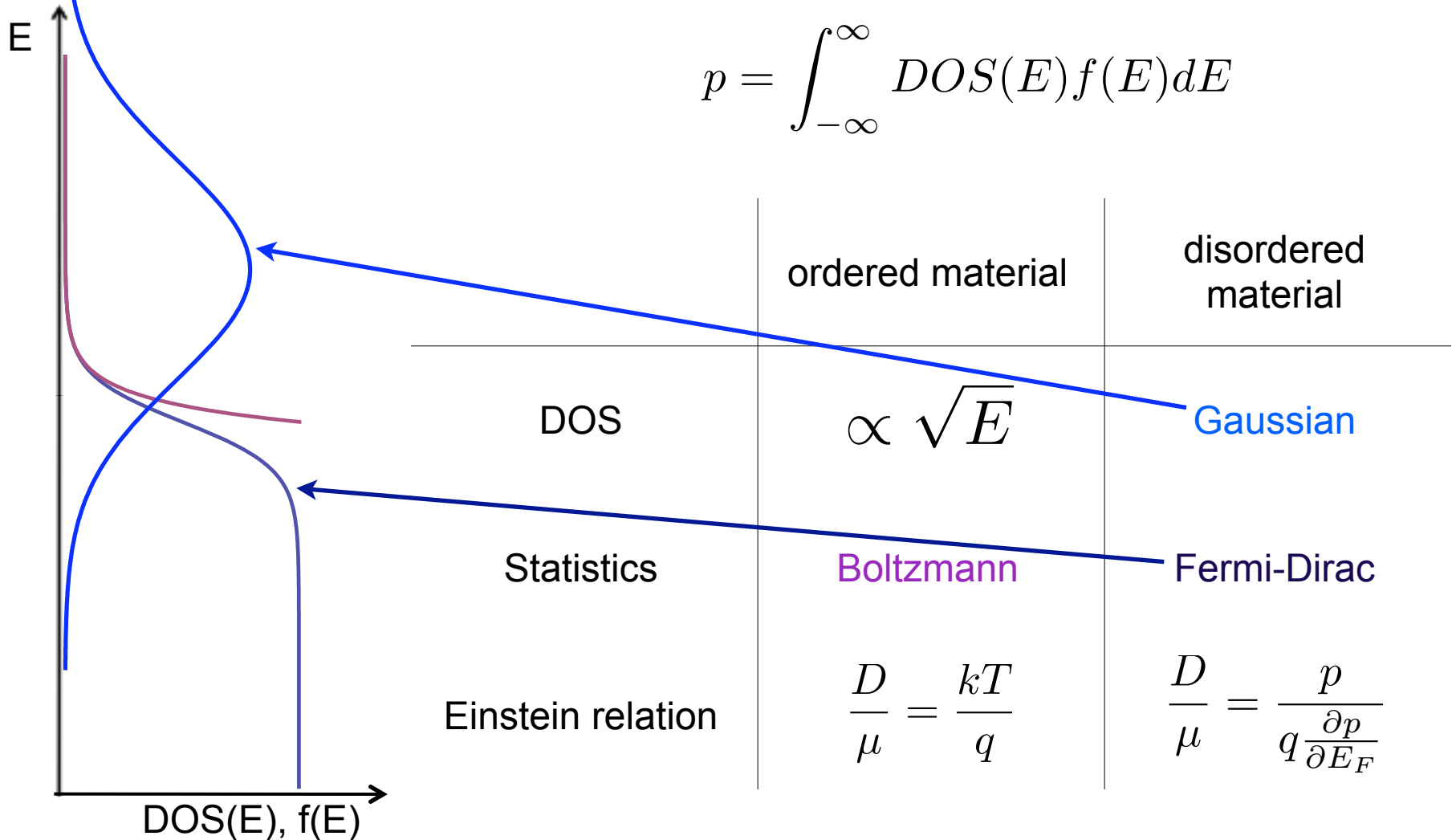
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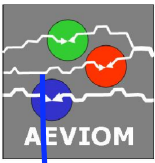

 mobility & diffusion coefficient
 are affected by the Gaussian DOS!



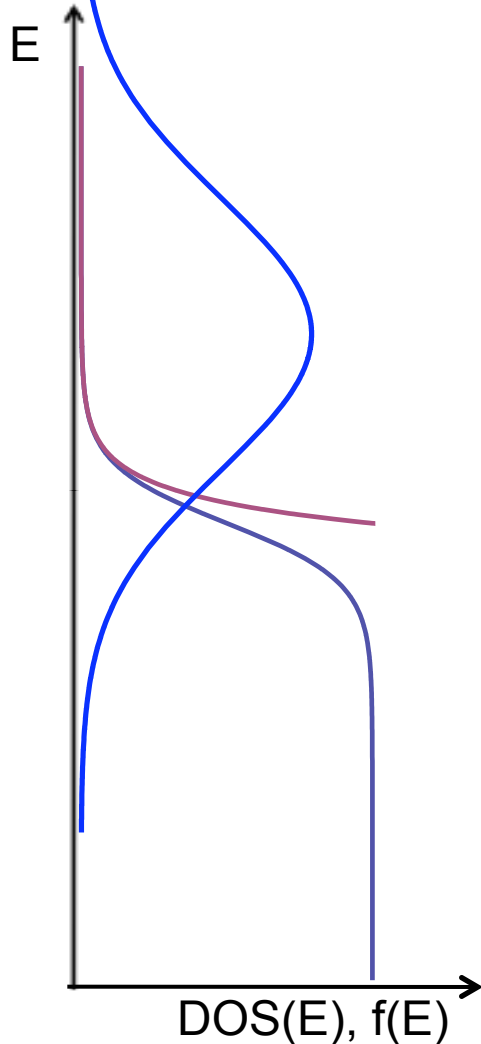


Generalized Einstein Relation





Generalized Einstein Relation

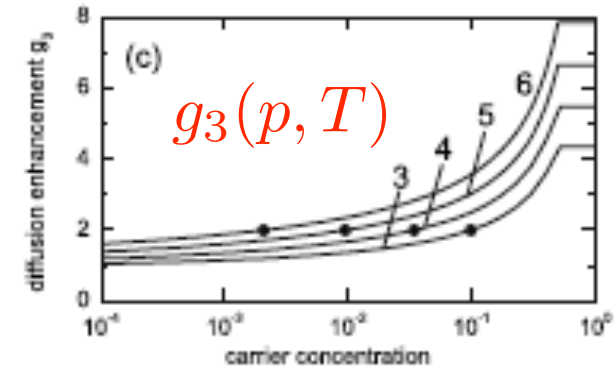
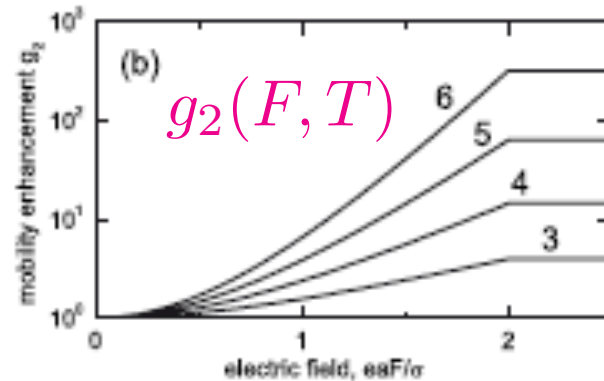
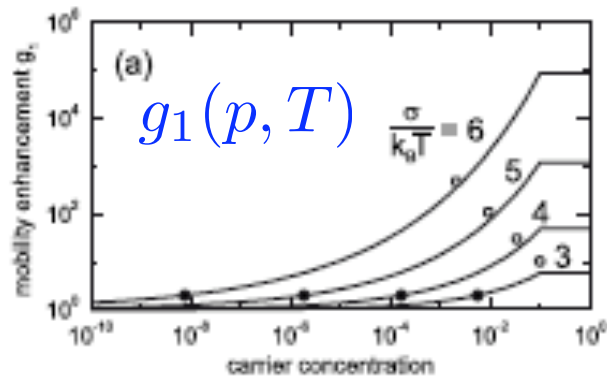


$$p = \int_{-\infty}^{\infty} DOS(E) f(E) dE$$

	ordered material	disordered material
DOS	$\propto \sqrt{E}$	Gaussian
Statistics	Boltzmann	Fermi-Dirac
Einstein relation	$\frac{D}{\mu} = \frac{kT}{q}$	$\frac{D}{\mu} = \frac{p}{q \frac{\partial p}{\partial E_F}}$

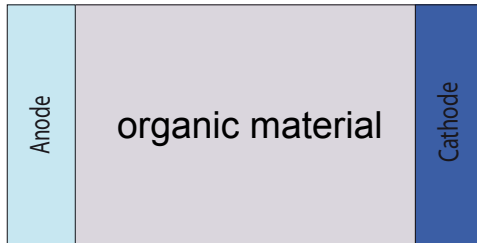
$$D_p = \frac{k_B T}{q} \mu_0(T, p, F) g_3(p, T)$$

$$\mu_p(T, p, F) = \mu_0(T) g_1(p, T) g_2(F, T)$$



Nonlinear equations for mobility and diffusion coefficient

Mobility depends on temperature, field and density

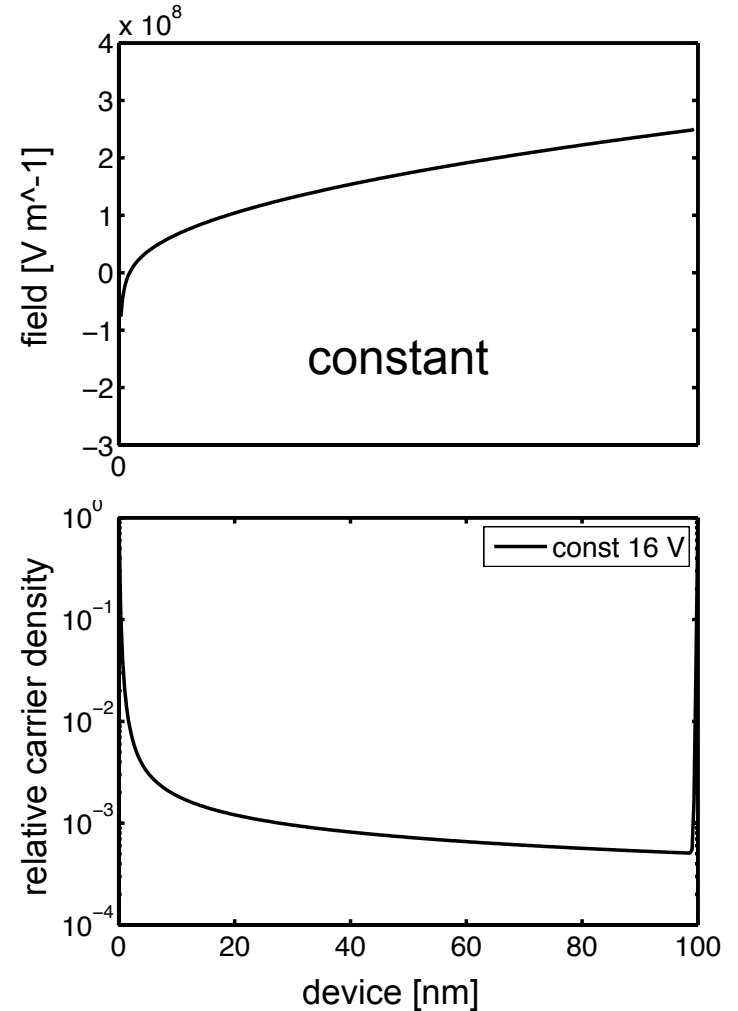


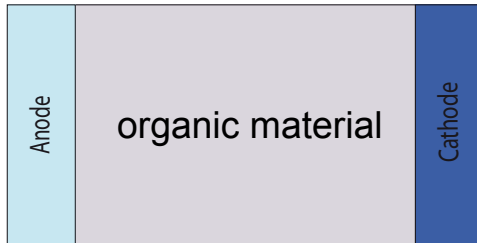
Assumption of ohmic contact:
Dirichlet boundary conditions

$$n_1 = 0.5N_t$$

$$n_2 = 0.5N_t$$

$$16V$$



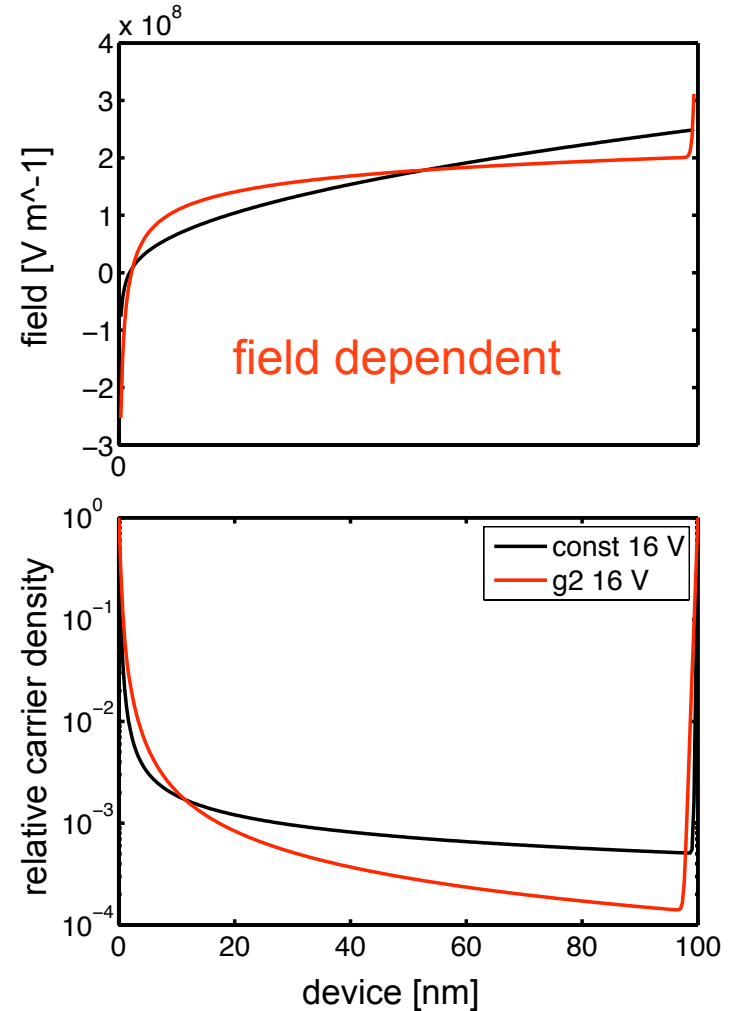


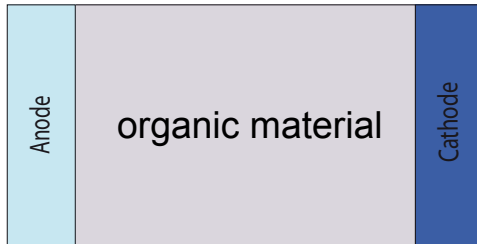
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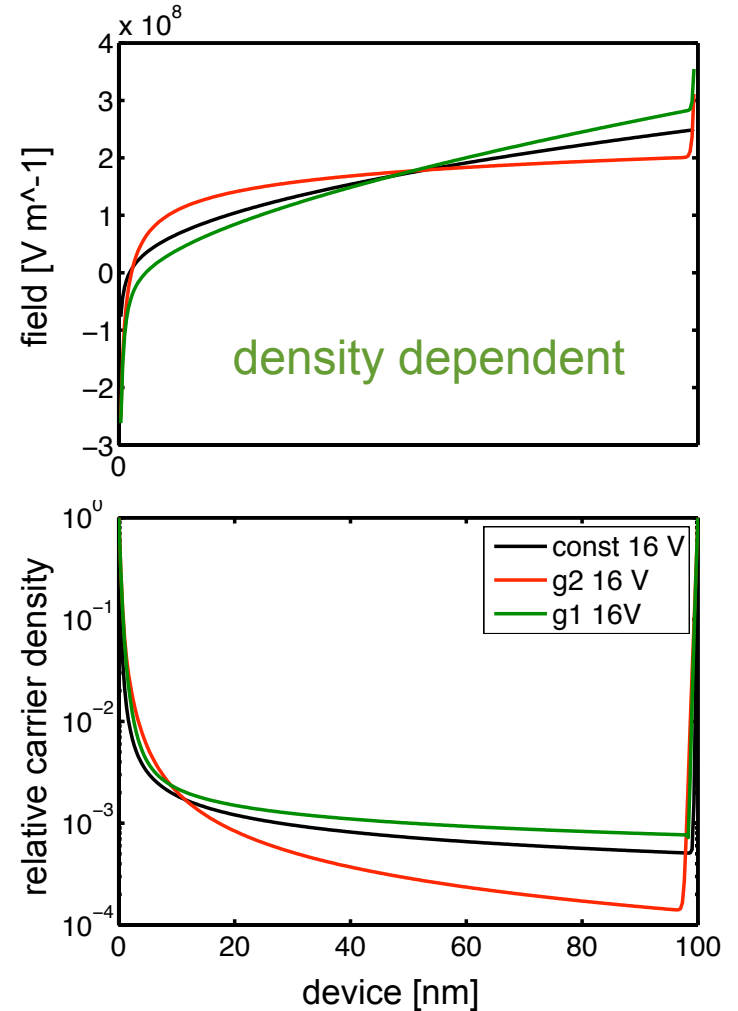


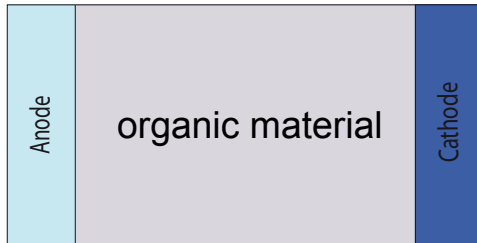
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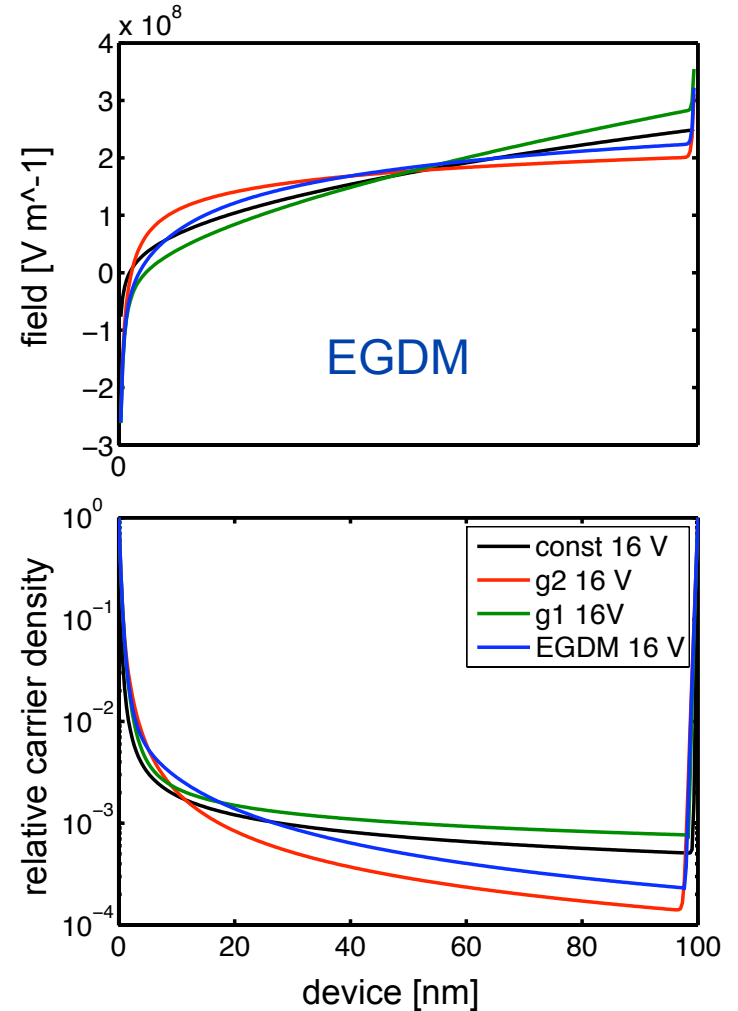


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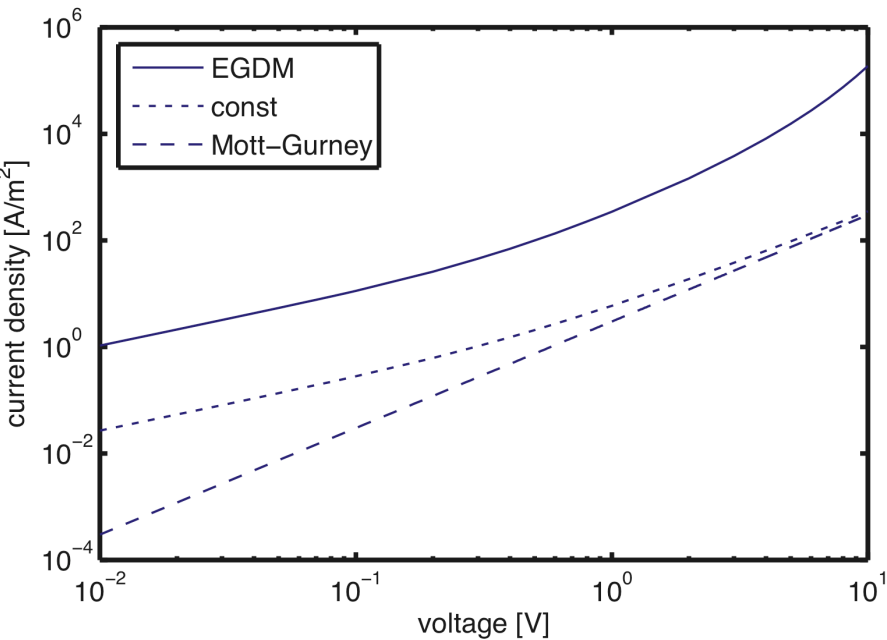
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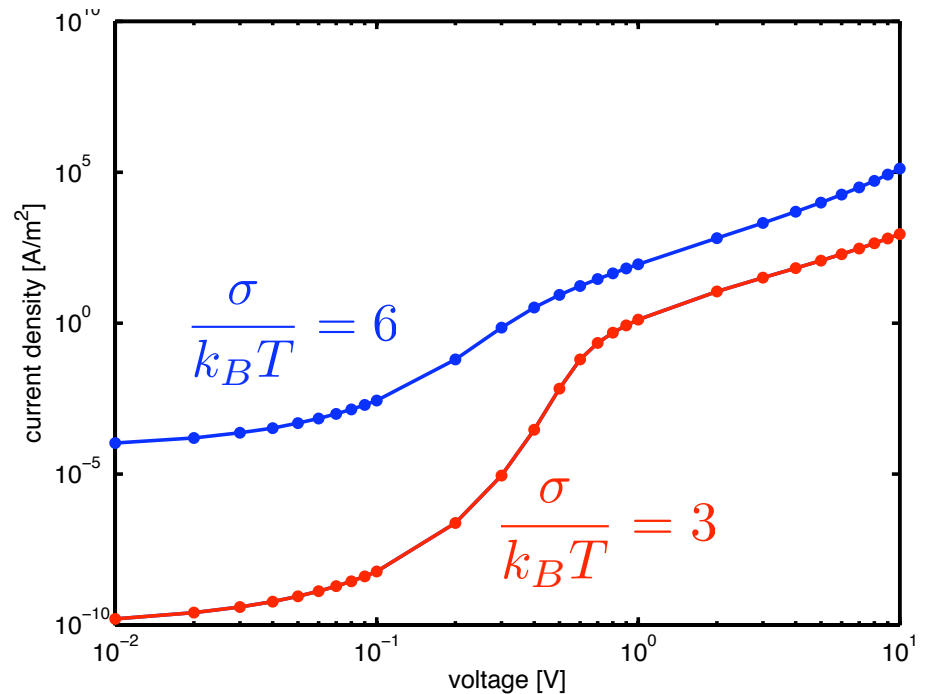


IV Curve
(hole-only device)



- Diffusion effects
- Field- and density-dependent

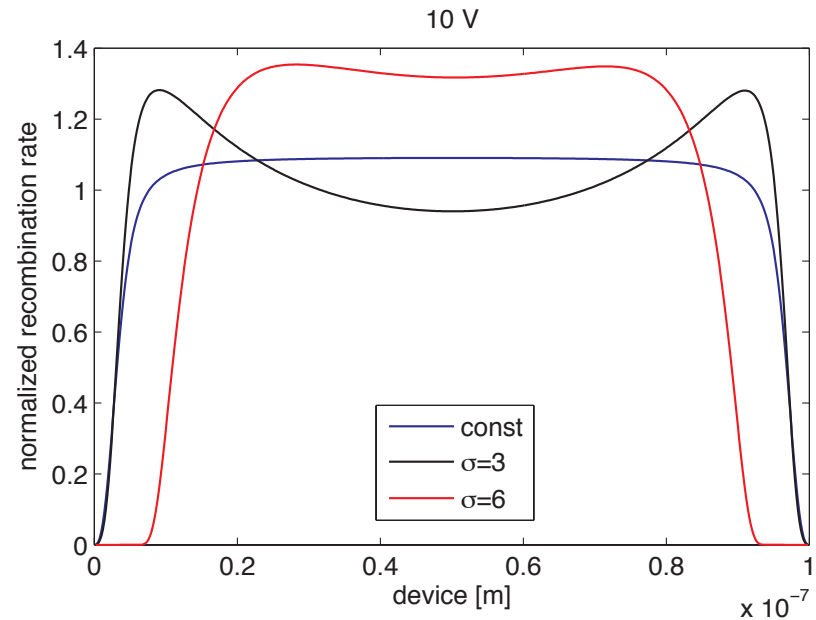
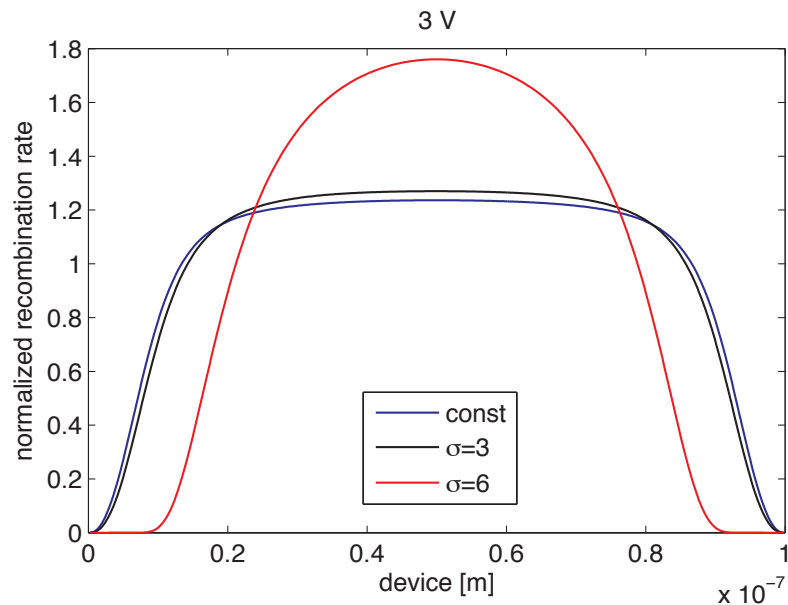
IV Curve
(hole-only device with 1eV built-in potential)



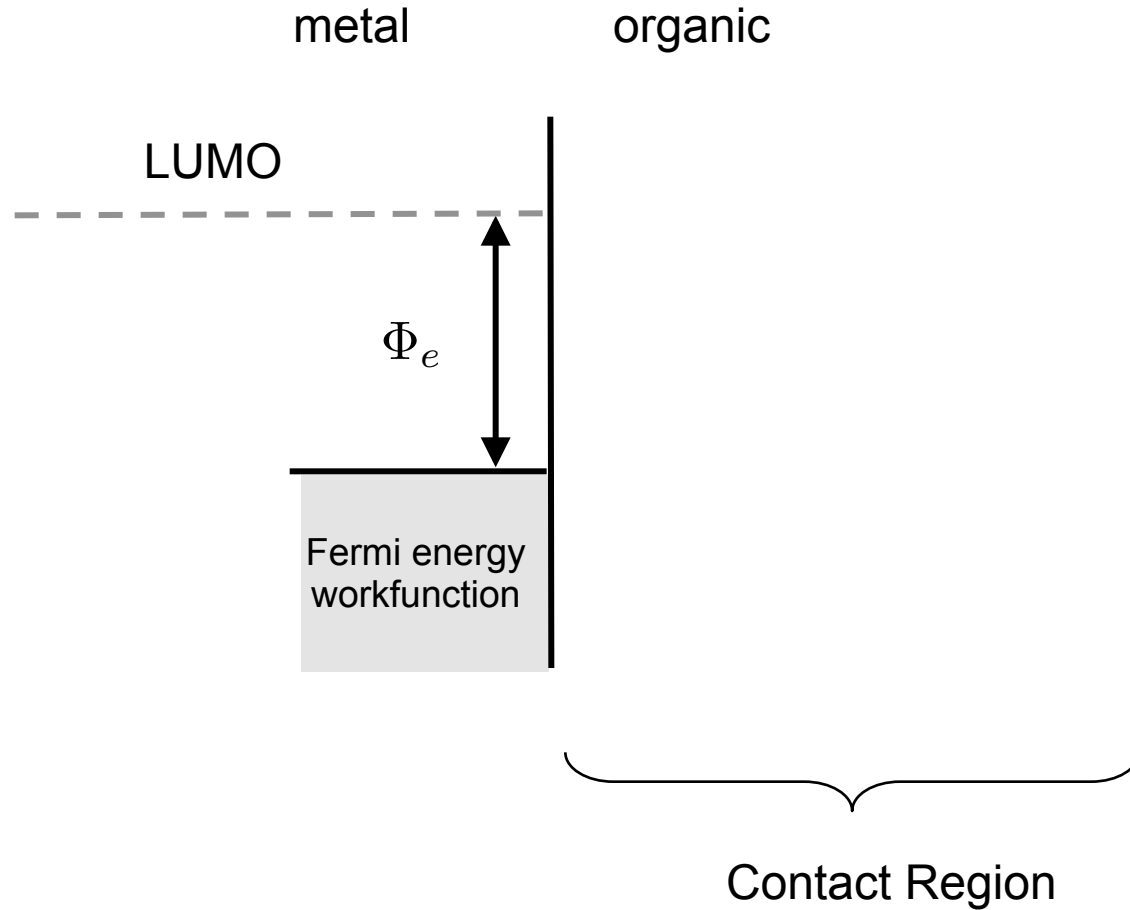
- Effects of different disorder parameters

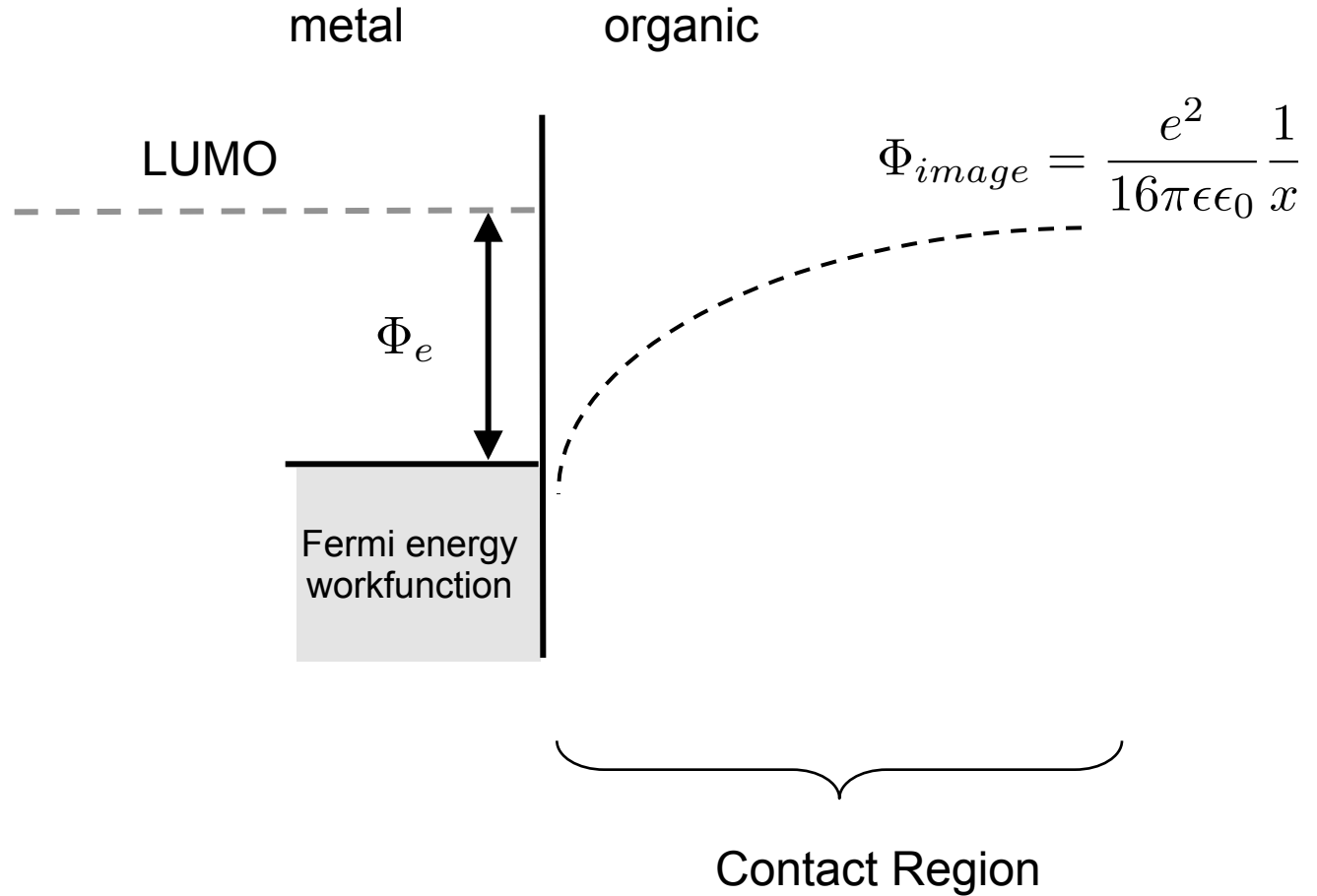
In good agreement with:

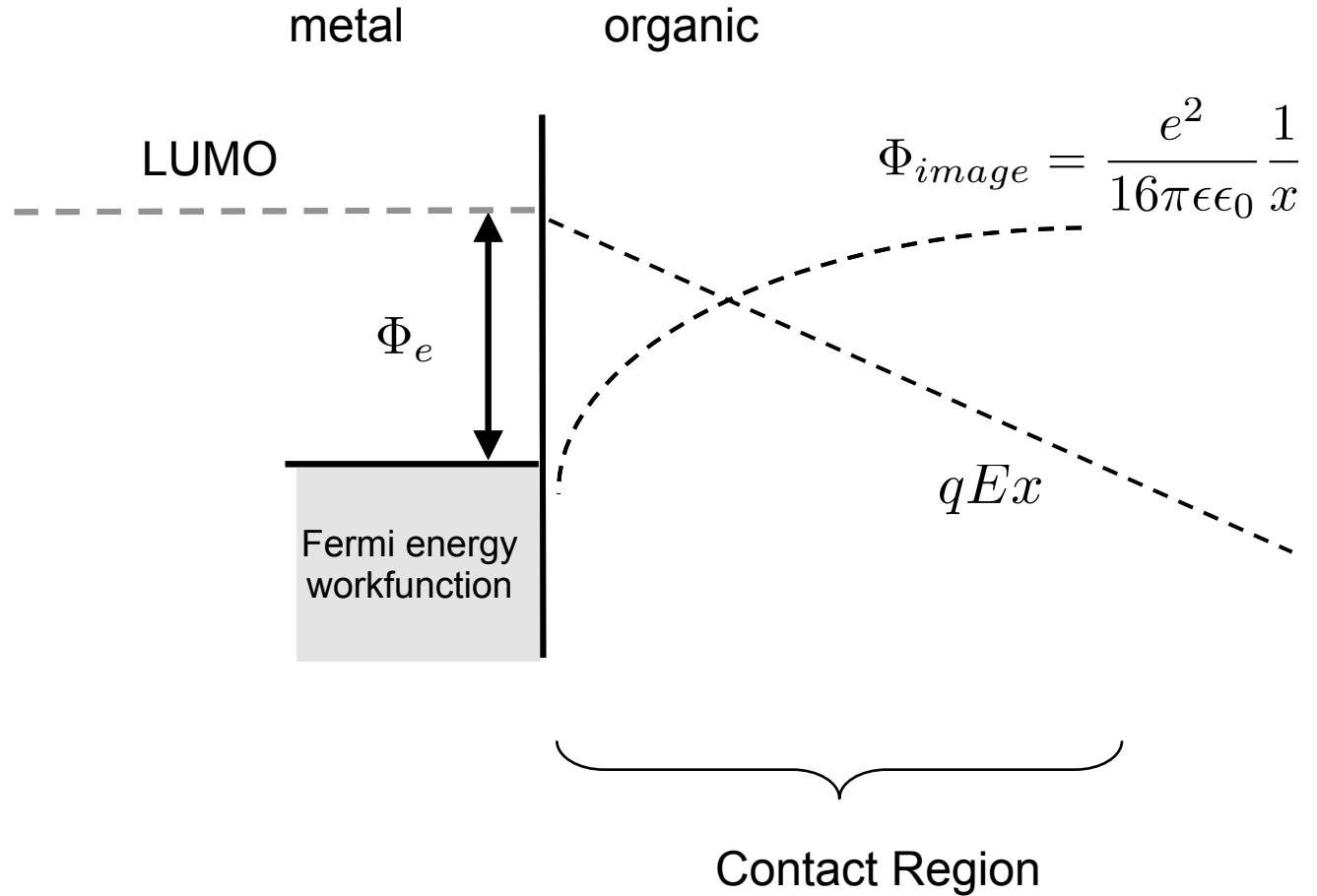
S. L. M. van Mensfoort, R. Coehoorn, Phys. Rev. B 78, 085207 (2008, Fig 9)

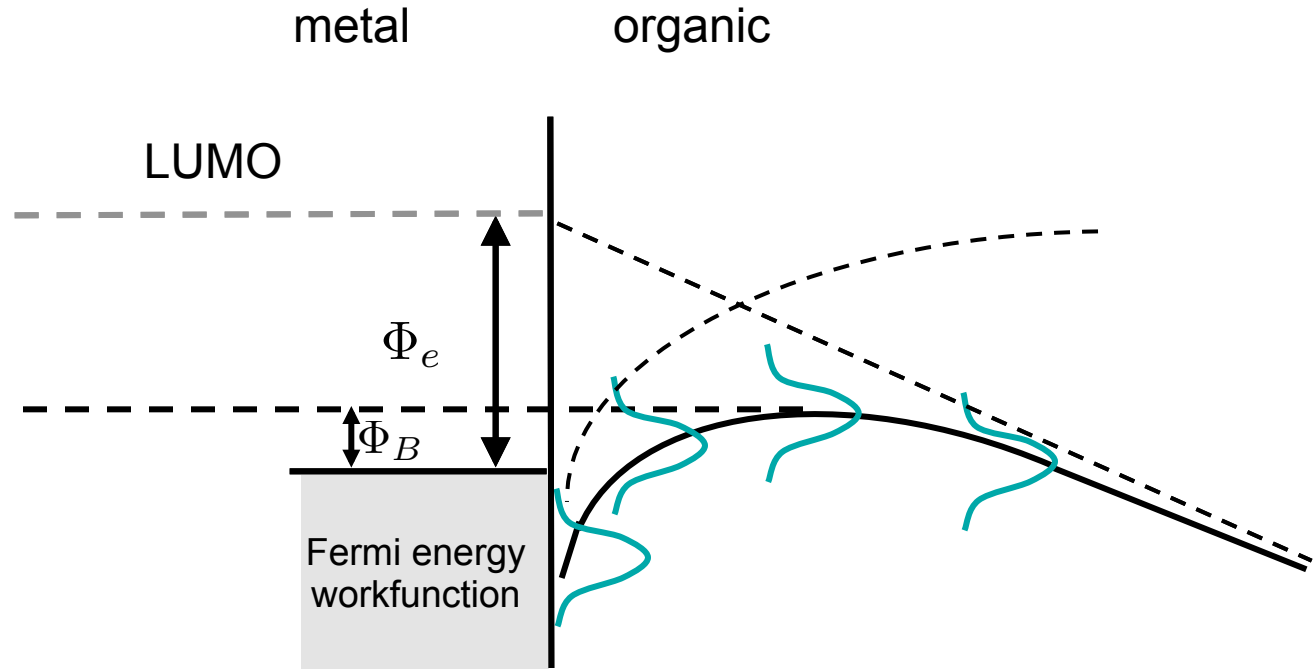


- Bipolar simulation with constant mobility and EGDM for $\hat{\sigma} = 3$ and $\hat{\sigma} = 6$
- Effects of disorder clearly visible

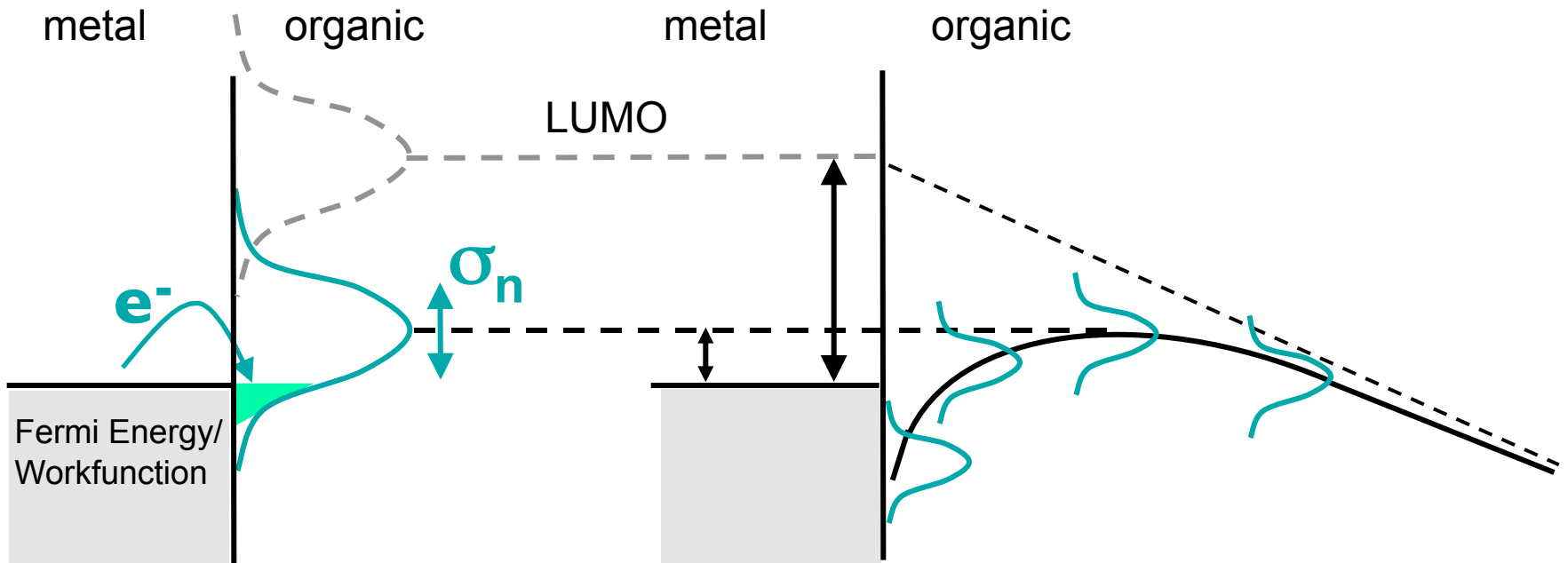








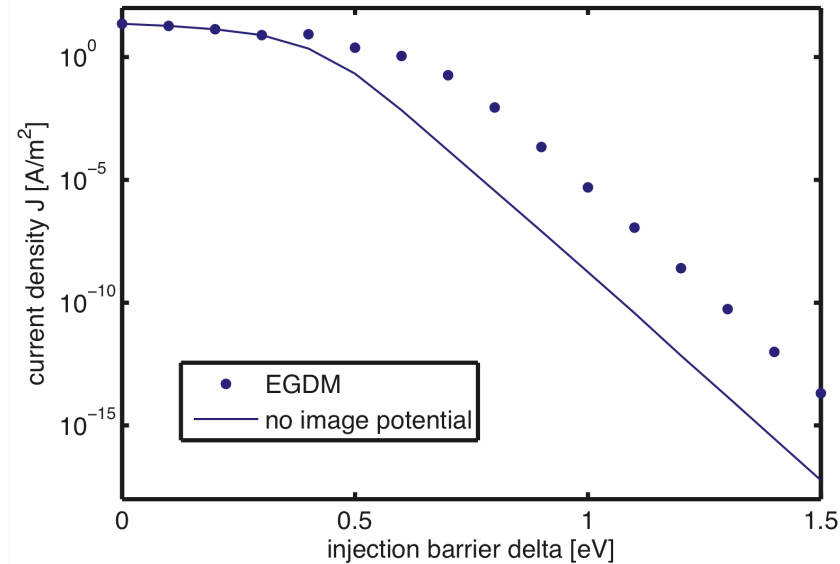
$$\Phi_e = eEx - \frac{e^2}{16\pi\epsilon\epsilon_0 x}$$



Density at contact depends on position of Gaussian DOS

Dependent boundary conditions

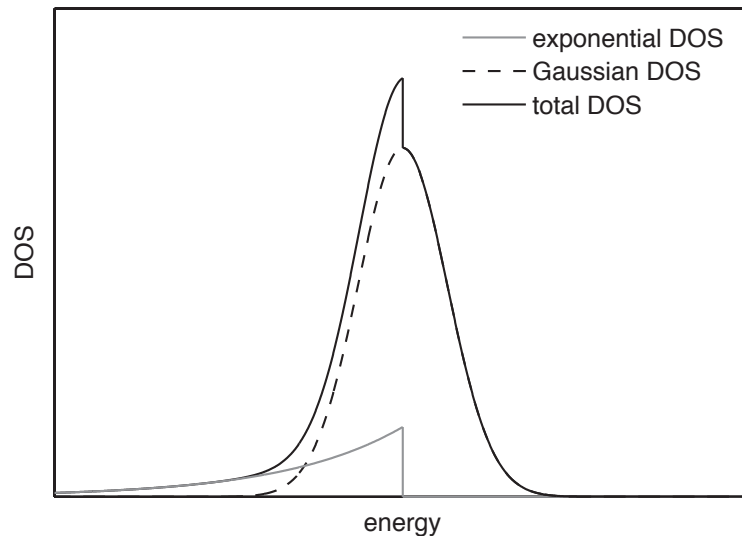
Dependence of the current density on the injection barrier at 2V



- No effect if injection barrier < 0.5 eV
- Higher currents with image potential
- Agrees with Monte Carlo results

In good agreement with:
 J.J.M. van der Holst, M.A. Uijtewaal, R. Balasubramanian, R. Coehoorn, P.A. Bobbert, G.A. de Wijs and R.A. de Groot (EUT, PRE), Phys. Rev. B (2009).

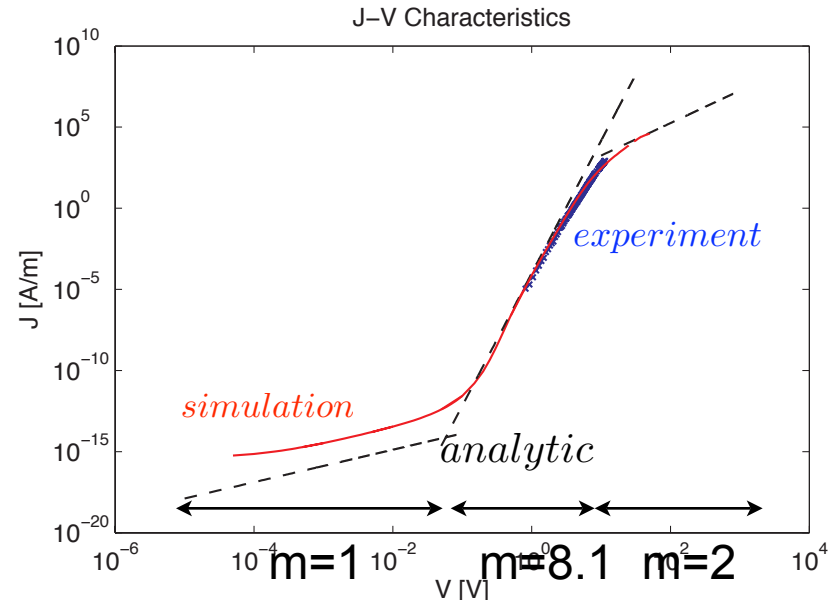
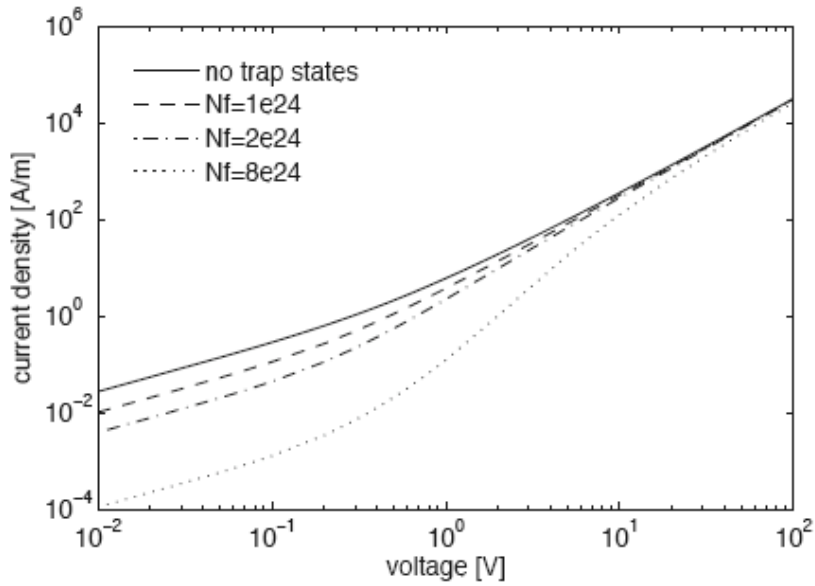
- localized sites with higher electron affinity
 - › impurities, chemical defects
- Model
 - › trap distribution: Exponential, Gaussian
 - › discrete levels: shallow, deep



$$\epsilon \Delta \psi = q(n - p + n_t - p_t)$$

$$\nabla \cdot J_p + q \frac{\partial p}{\partial t} = -qR(p, n)$$

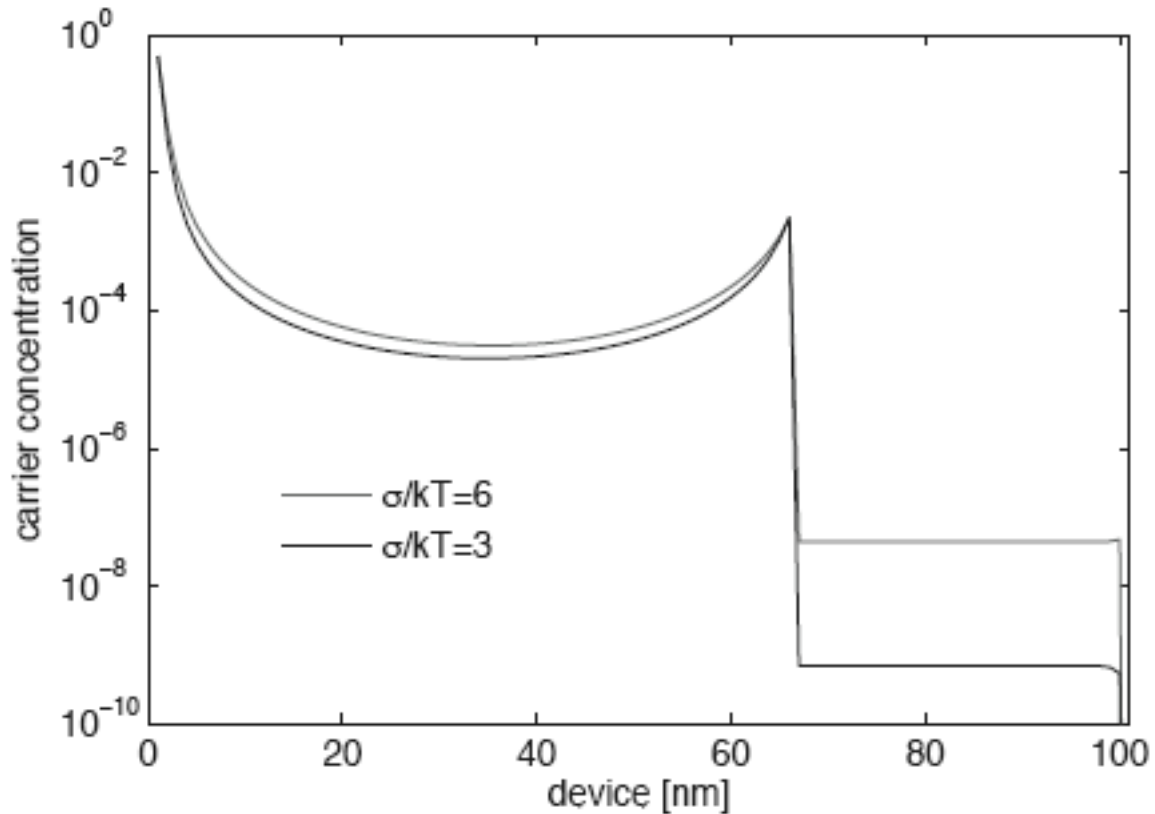
$$J_p = -q\mu_p p \nabla \psi - qD_p \nabla p$$



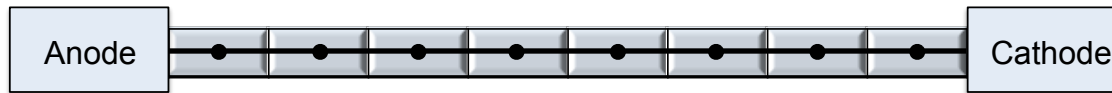
trap density influences current density

Analytical solution for Gaussian DOS:
M. M. Mandoc, B. de Boer, G. Paasch, P. W. M. Blom, Phys. Rev. B (2007).

- Stack of organic material to optimize recombination profiles and light emission



- 1-dimensional finite volume method
 - › Domain divided into n grid points



- Reformulation of problem

$$F_1(\psi, p, n) = \epsilon \Delta \psi - q(n - p) \stackrel{!}{=} 0$$

$$F_2(\psi, p, n) = \nabla \cdot (-q\mu_p p \nabla \psi - qD_p \nabla p) + q \frac{\partial p}{\partial t} + qR \stackrel{!}{=} 0$$

$$F_3(\psi, p, n) = \nabla \cdot (-q\mu_n n \nabla \psi + qD_n \nabla n) - q \frac{\partial n}{\partial t} - qR \stackrel{!}{=} 0$$

- Integration over each box

- Neglecting recombination and assuming a constant current density through the device

$$q\mu_n(U_t \frac{\partial n}{\partial x} - n \frac{\partial \psi}{\partial x}) = c$$

- Boundary values $n(x_{i-1}) = n_{i-1}$ and $n(x_i) = n_i$
- Analytic solution

$$n(x) = n_{i-1}(1 - g(x)) + n_i g(x)$$

with

$$g(x) = \frac{1 - \exp\left(\frac{(\psi_i - \psi_{i-1})}{U_t} \frac{x - x_{i-1}}{x_i - x_{i-1}}\right)}{1 - \exp\left(\frac{\psi_i - \psi_{i-1}}{U_t}\right)}$$

- Analytic solution serves as Ansatz function
 - › Scharfetter-Gummel discretization

- Exponential fitting for drift-diffusion (F2 and F3)
 - › Scharfetter-Gummel discretization with generalized Einstein relation and density- and fielddependent mobility

- System of (3 x n) strongly coupled equations

$$\vec{F}(\vec{x}) = \begin{pmatrix} \vec{F}_1(\vec{x}) \\ \vec{F}_2(\vec{x}) \\ \vec{F}_3(\vec{x}) \end{pmatrix} \quad \vec{x} = \begin{pmatrix} \psi_1 \\ \vdots \\ \psi_n \\ n_1 \\ \vdots \\ n_n \\ p_1 \\ \vdots \\ p_n \end{pmatrix}$$

- Dirichlet boundary conditions:
 - › Values for potential and carriers given at electrodes

- Variables sets

- › carrier concentrations (ψ, p, n)

- › quasi-Fermi level (ψ, ϕ_p, ϕ_n)

- Assumption: Boltzmann statistics

$$p = n_{int,eff} \exp\left(\frac{q(\phi_p - \psi)}{kT}\right)$$

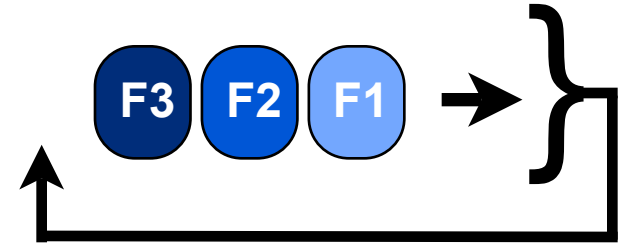
$$n = n_{int,eff} \exp\left(\frac{q(\psi - \phi_n)}{kT}\right)$$

- › Slotboom (ψ, Φ_p, Φ_n)

$$\Phi_p = \exp\left(\frac{q\phi_p}{kT}\right) \quad p = p_i \Phi_p \exp\left(\frac{-q\psi}{kT}\right)$$

$$\Phi_n = \exp\left(\frac{-q\phi_n}{kT}\right) \quad n = n_i \Phi_n \exp\left(\frac{q\psi}{kT}\right)$$

- De-coupled solving
 - › Gummel algorithm
- Coupled solving
 - › Newton algorithm



Find x^* so that $F(x^*)=0$.

$$F(x) = F(x^*) + J(x^*)(x - x^*)$$

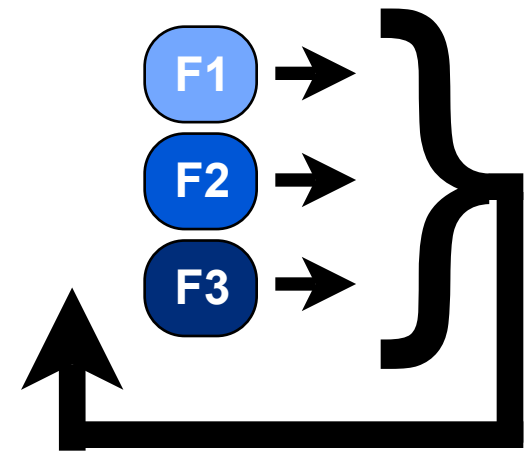
$$J(x) = \begin{bmatrix} \frac{\partial F_1(x)}{\partial x_1} & \dots & \frac{\partial F_1(x)}{\partial x_N} \\ \vdots & \ddots & \vdots \\ \frac{\partial F_N(x)}{\partial x_1} & \dots & \frac{\partial F_N(x)}{\partial x_N} \end{bmatrix}$$

$$\Rightarrow x^{k+1} = x^k - J(x^k)^{-1} F(x^k)$$

Taylor Series

Jacobian Matrix

Iteration function



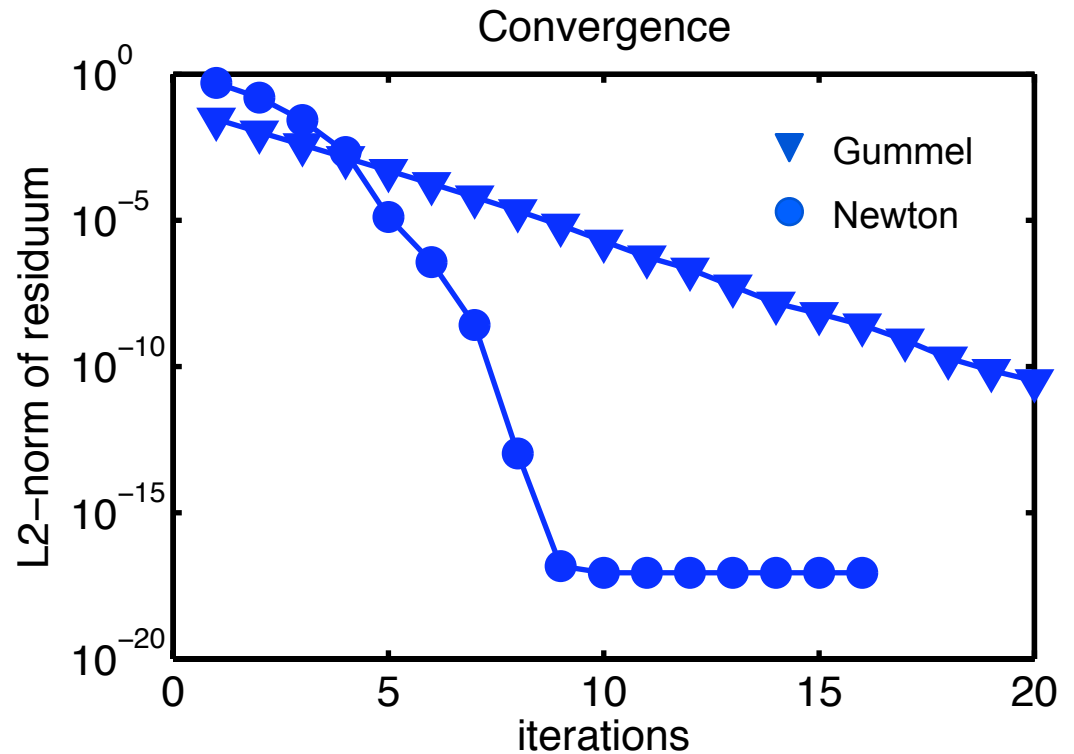
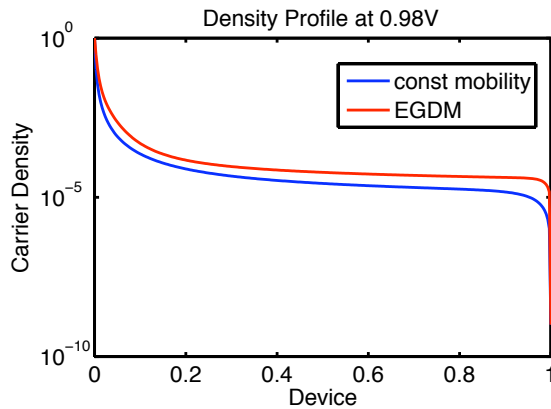


Algorithms

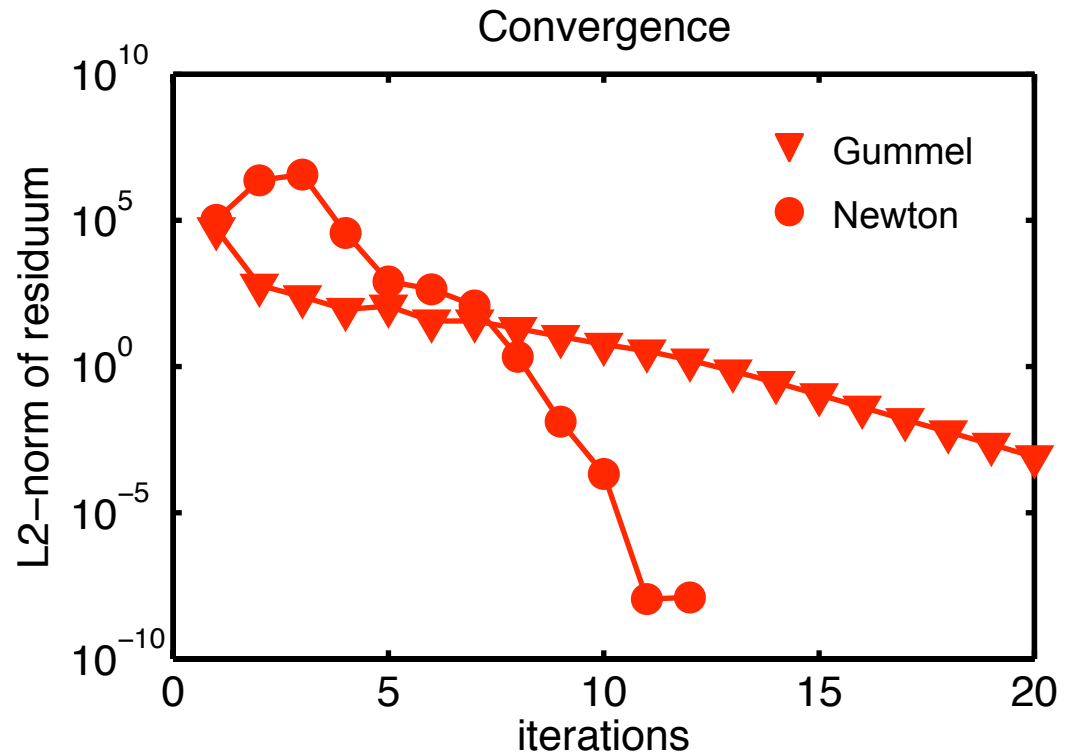
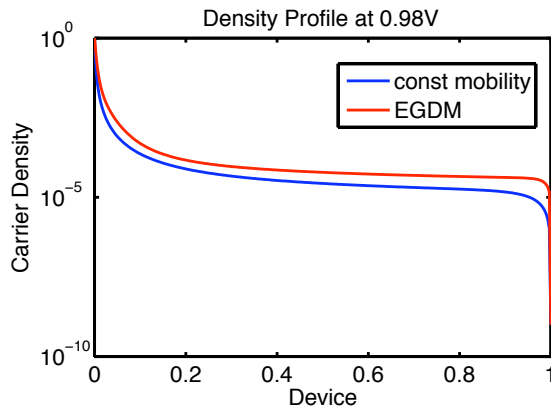


- Gummel
 - › steady-state
 - › transient
- Newton
 - › steady-state
 - › transient
- Initial guess
 - › no bias applied, Boltzmann approximation
- Gummel steady-state
 - › Damping
- Newton
 - › Damping
 - › Homotopy

$$\text{L2-Norm: } |F| = \sqrt{\sum_{k=1}^n |F_k|^2}$$

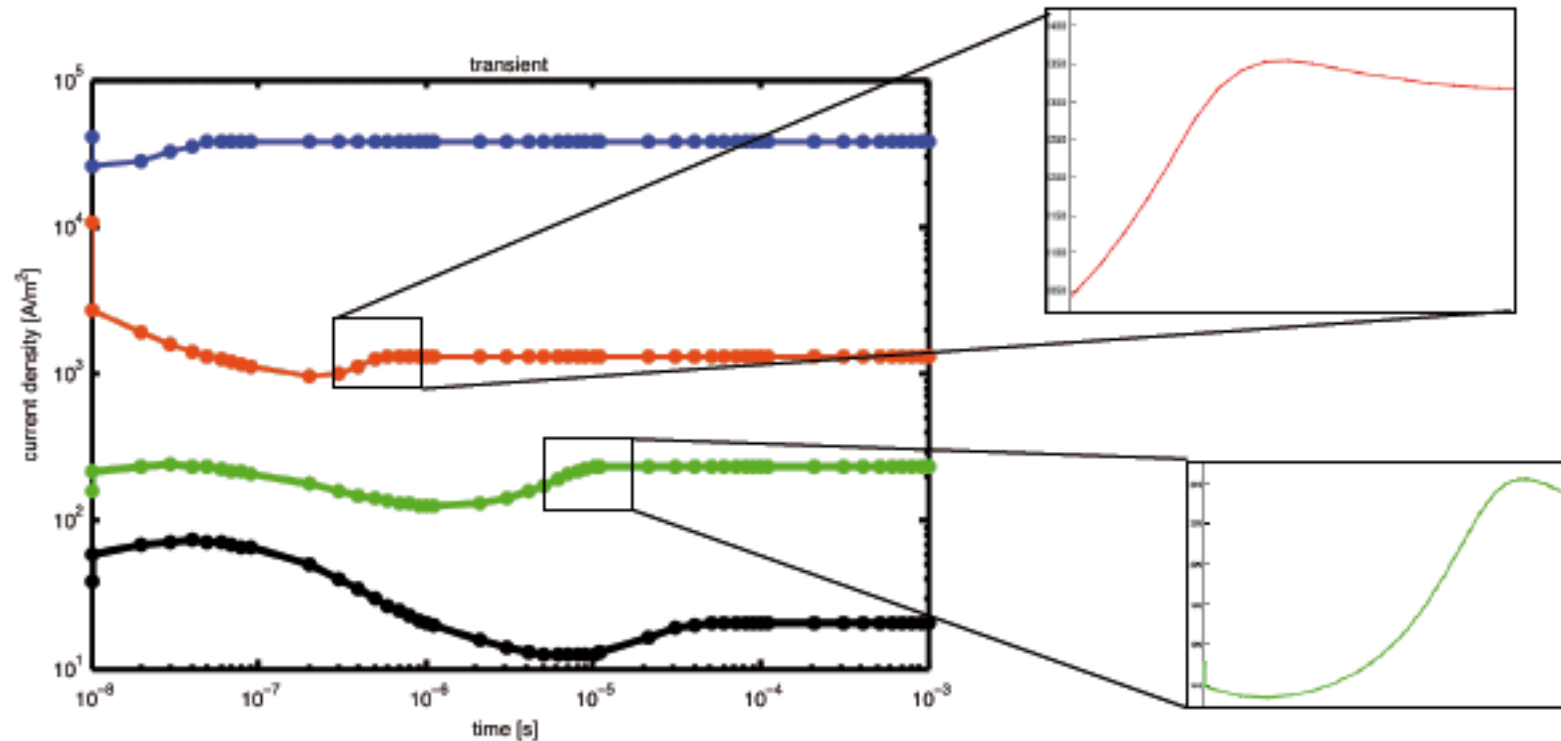


$$\text{L2-Norm: } |F| = \sqrt{\sum_{k=1}^n |F_k|^2}$$



- Convergence for Gummel and Newton algorithm
- Fewer iterations needed for Newton algorithm

- Implicit Euler time step

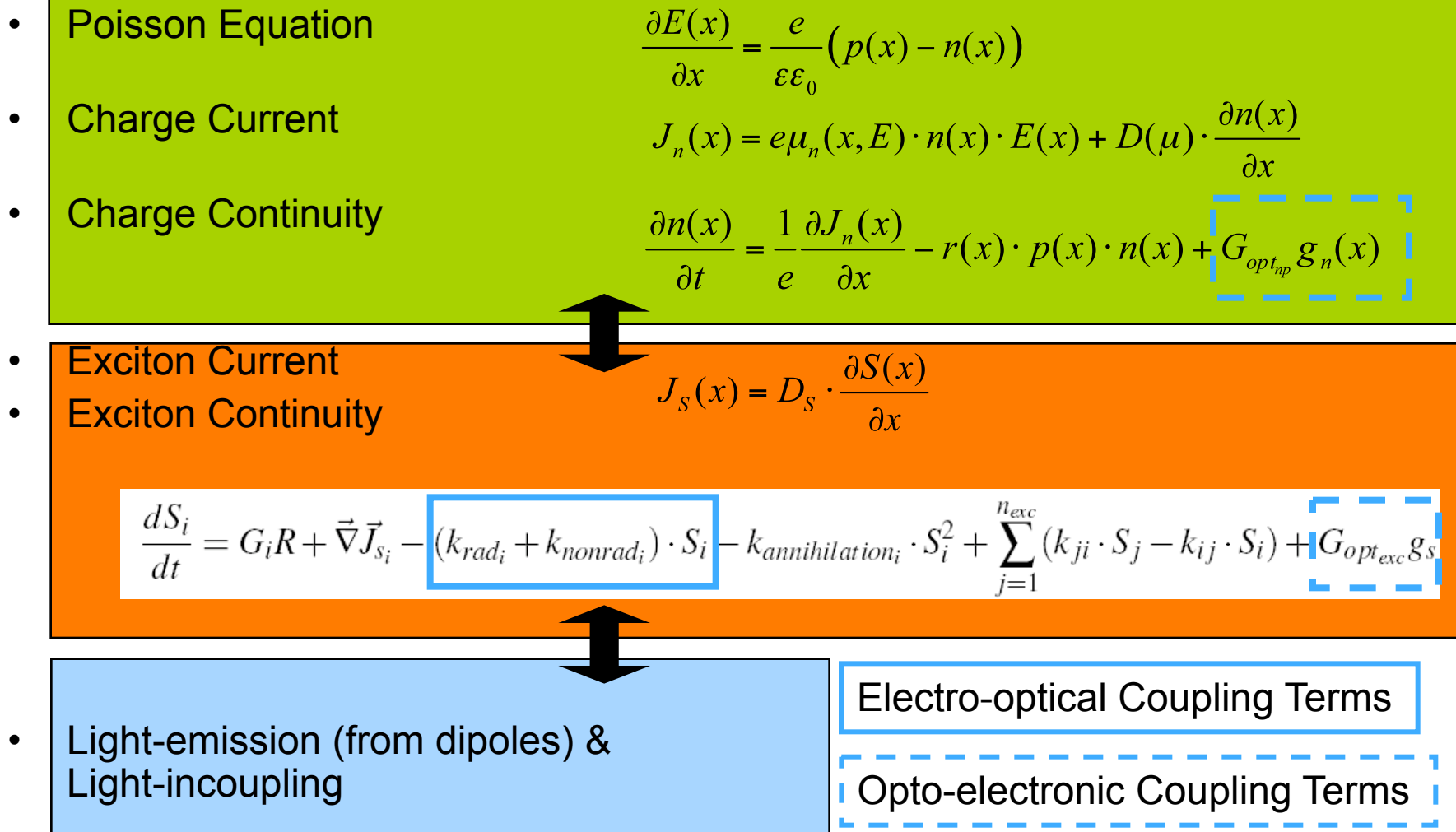




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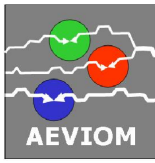
Extended version of the models published by Ruhstaller et al., J. Appl. Phys. **89**, 4575, (2001) and Ruhstaller et al., IEEE JSTQE **9**, (3) 723, (2003)



Outlook



- ✓ Modeling of charge carrier transport (1st generation)
 - › Gummel
 - › Newton
- ✓ Bipolar (1st generation)
- ✓ Injection (2nd generation)
- ✓ Organic material properties
 - › Disorder (2nd generation)
 - › Mobility (2n generation)
 - › Generalized Einstein relation (2nd generation)
- ✓ Traps (2nd generation)
- ✓ Multilayer OLEDs (1st generation)
 - Exciton dynamics (1st generation)
 - Parameter extraction
 - Optical simulations
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Acknowledgement



- We acknowledge the financial support of RF7
- Thanks for your attention!